THE NUCLEAR OVERHAUSER EFFECT IN STRUCTURAL AND CONFORMATIONAL ANALYSIS

Second Edition

Methods in Stereochemical Analysis

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THE NUCLEAR OVERHAUSER EFFECT IN STRUCTURAL AND CONFORMATIONAL ANALYSIS

Second Edition

David Neuhaus Michael P. Williamson



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CONTENTS

DDE	FACE	W
PKE	EFACE	XV
ACk	KNOWLEDGMENTS	xvii
PRE	FACE TO THE FIRST EDITION	xix
SYN	MBOLS, ABBREVIATIONS, AND UNITS	xxi
INTI	RODUCTION	XXV
PAR	RT I. THEORY	2
СНА	PTER 1. BACKGROUND	3
1.1.	Energy Levels, Populations, and Intensities / 3	
1.2.	Relaxation, T_1 and $T_2 / 8$	
1.3.	The Nature of Relaxation / 12	
1.4.	The Local Field and Dipole-Dipole Relaxation / 13	
1.5.	Pulses and Saturation / 15	
	References / 22	

CHAPTER 2.		THE STEADY-STATE NOE FOR TWO SPINS		
2.1.	The Orig	gin and Form of the NOE / 23		
	2.1.1.	Qualitative Considerations / 23		
	2.1.2.	The Solomon Equations / 27		
2.2.	Depende	ence of the NOE on Molecular Motion / 30		
	2.2.1.	Correlation Times, Spectral Density Functions, and Transition Probabilities / 31		
	2.2.2	Anisotropic Tumbling / 38		
2.3	What the	e Symbols Mean for Two Spins and for Many Spins / 38		
	2.3.1.	Relaxation Rates / 39		
	2.3.2.	T_1 Measurements and Cross-Relaxation / 42		
2.4.	Effects of	of Other Relaxation Sources / 46		
	2.4.1.	The External Relaxation Rate ρ^* / 46		
	2.4.2.	Intermolecular Dipole-Dipole Relaxation / 50		
	2.4.3.	Quadrupolar Relaxation / 52		
	2.4.4.	Chemical Shift Anisotropy (CSA) Relaxation / 53		
	2.4.5.	Scalar Relaxation / 53		
	2.4.6.	Spin-Rotation Relaxation / 53		
2.5.	The Hete	eronuclear NOE / 54		
2.6.	An Exter	nsion to the Solomon Equations / 59		
	Reference	ees / 60		
		THE STEADY-STATE NOE IN RIGID	62	
3.1.	The Equ	ations / 63		
	3.1.1.	The Solomon Equations for More Than Two Spins / 63		
	3.1.2.	Cross-Correlation / 66		
	3.1.3.	Two General Solutions to the Multispin Solomon Equations / 68		
		3.1.3.1. Saturation of One Spin / 68		
		3.1.3.2. Saturation of All Spins Except One / 69		
	3.1.4.	Internuclear Distances and Steady-State NOE Enhancements / 71		
3.2.	What the	e Equations Mean / 72		
	3.2.1.	General: Direct Enhancements and Spin Diffusion / 72		
	3.2.2.	Interpretation at the Extreme Narrowing Limit $(\omega \tau_c \ll 1) / 76$		

		3.2.2.2. Indirect Effects / 79	
		3.2.2.3. When Do Indirect Effects Matter? / 80	
		3.2.2.4. Magnetic Equivalence / 82	
		3.2.2.5. T_1 and the 3/2 Effect / 84	
		3.2.2.6. Chemical Equivalence / 85	
	3.2.3.	Away from the Extreme Narrowing Limit / 86	
3.3.	In Practi	ice / 91	
	3.3.1.	Incomplete Saturation / 92	
	3.3.2.	Failure to Reach Steady-State / 93	
	3.3.3.	Competition From Other Relaxation Sources / 94	
		References / 97	
CHAF	PTER 4.	THE KINETICS OF THE NOE	98
4.1.	Introduc	tion / 98	
7.1.	4.1.1.	Types of Kinetic NOE Experiment / 99	
	4.1.2.	Overview / 100	
4.2.		of the Kinetic NOE in a Two-Spin System / 105	
4.3.		of the Kinetic NOE in Multispin Systems / 108	
	4.3.1.	Multispin Kinetics in Transient NOE Experiments / 108	
	4.3.2.	Multispin Kinetics in TOE Experiments / 111	
4.4.	Estimati	ng Internuclear Distances / 111	
	4.4.1.	The Initial Rate Approximation / 112	
	4.4.2.	Distances From Enhancement Ratios / 113	
	4.4.3.	Errors in Distance Measurements Using the Initial Rate Approximation / 115	
	4.4.4.	Spin Diffusion, Nonlinear NOE Growth, and	
	7.7.7.	Interpretation / 117	
4.5.	More Al	pout Experiments / 122	
	4.5.1.	Symmetry in Kinetic NOE Experiments / 122	
	4.5.2.	T_1 Values as an Aid to Interpretation / 124	
	Reference	ces / 127	
CHAF	PTER 5.	THE EFFECTS OF EXCHANGE AND	
_	RNAL MO		129
5.1.	Transfer	of Saturation / 131	
5.2.	General	Equations for the NOE in Systems of Two-Site	

Exchange / 136

3.2.2.1. Direct Effects / 76

	5.2.1.	Exchange in a Two-Spin System / 136	
	5.2.2.	Exchange in Dimethylformamide / 143	
5.3.		tions to More Complicated Cases of Exchange / 148	
	5.3.1.	Averaging of Rates Rather than Enhancements / 148	
	5.3.2.	Analyzing Conformational Equilibria / 150	
		5.3.2.1. Olefinic Methoxy Conformations / 150	
		5.3.2.2. Nucleotide Conformations / 151	
		5.3.2.3. A Statistical Approach / 156	
5.4.	Estimat Analysi	ing Flexibility Using Heteronuclear Relaxation s / 158	
5.5	How In	ternal Motions Average Internuclear Distances / 167	
	5.5.1.	Internal Motions Slower Than Overall Tumbling: r^{-6} Averaging" / 171	
	5.5.2.	Internal Motions Faster Than Overall Tumbling / 172	
5.6.	Allowir	ng for Averaging / 174	
5.7.	The Tra	unsferred NOE / 178	
5.8.	Intermo	lecular NOE Enhancements Involving Water / 185	
		ces / 187	
CHAF	PTER 6.	COMPLICATIONS FROM SPIN-SPIN COUPLING	190
6.1.	Decoup	ling / 190	
	-	ling / 190 re Population Transfer / 191	
6.1. 6.2.	-	re Population Transfer / 191	
	Selectiv	re Population Transfer / 191 Theory / 192	
	Selective 6.2.1. 6.2.2.	re Population Transfer / 191 Theory / 192 Consequences / 199	
6.2.	Selective 6.2.1. 6.2.2. Effects	re Population Transfer / 191 Theory / 192 Consequences / 199 of Cross-Correlation / 201	
6.2.6.3.	Selective 6.2.1. 6.2.2. Effects Strong	re Population Transfer / 191 Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204	
6.2.6.3.	Selective 6.2.1. 6.2.2. Effects Strong 6.4.1.	re Population Transfer / 191 Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204 A{B} Enhancements / 206	
6.2.6.3.	Selective 6.2.1. 6.2.2. Effects Strong 6.4.1. 6.4.2.	Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204 $A\{B\} \text{ Enhancements } / 206$	
6.2.6.3.	Selective 6.2.1. 6.2.2. Effects Strong 6.4.1.	Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204 A{B} Enhancements / 206 AB{X} Enhancements / 206 Scalar Relaxation / 213	
6.2.6.3.	Selective 6.2.1. 6.2.2. Effects Strong 6.4.1. 6.4.2.	Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204 A{B} Enhancements / 206 AB{X} Enhancements / 206 Scalar Relaxation / 213 6.4.3.1. Scalar Relaxation of the First Kind / 213	
6.2.6.3.	Selective 6.2.1. 6.2.2. Effects Strong 6.4.1. 6.4.2. 6.4.3.	Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204 A{B} Enhancements / 206 AB{X} Enhancements / 206 Scalar Relaxation / 213	
6.2. 6.3. 6.4.	Selective 6.2.1. 6.2.2. Effects Strong 6.4.1. 6.4.2. 6.4.3.	Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204 A{B} Enhancements / 206 AB{X} Enhancements / 206 Scalar Relaxation / 213 6.4.3.1. Scalar Relaxation of the First Kind / 213 6.4.3.2. Scalar Relaxation of the Second Kind / 216 ces / 217	210
6.2. 6.3. 6.4.	Selective 6.2.1. 6.2.2. Effects Strong 6.4.1. 6.4.2. 6.4.3.	Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204 A{B} Enhancements / 206 AB{X} Enhancements / 206 Scalar Relaxation / 213 6.4.3.1. Scalar Relaxation of the First Kind / 213 6.4.3.2. Scalar Relaxation of the Second Kind / 216 EXPERIMENTAL	219
6.2. 6.3. 6.4.	Selective 6.2.1. 6.2.2. Effects Strong 6.4.1. 6.4.2. 6.4.3.	Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204 A{B} Enhancements / 206 AB{X} Enhancements / 206 Scalar Relaxation / 213 6.4.3.1. Scalar Relaxation of the First Kind / 213 6.4.3.2. Scalar Relaxation of the Second Kind / 216 ces / 217	219 22 1
6.2. 6.3. 6.4.	Selective 6.2.1. 6.2.2. Effects Strong 6.4.1. 6.4.2. 6.4.3. Referen	Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204 A{B} Enhancements / 206 AB{X} Enhancements / 206 Scalar Relaxation / 213 6.4.3.1. Scalar Relaxation of the First Kind / 213 6.4.3.2. Scalar Relaxation of the Second Kind / 216 EXPERIMENTAL ONE-DIMENSIONAL EXPERIMENTS Preparation / 221	
6.2. 6.3. 6.4. PAR CHAF	Selective 6.2.1. 6.2.2. Effects Strong 6.4.1. 6.4.2. 6.4.3. Referen	Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204 A{B} Enhancements / 206 AB{X} Enhancements / 206 Scalar Relaxation / 213 6.4.3.1. Scalar Relaxation of the First Kind / 213 6.4.3.2. Scalar Relaxation of the Second Kind / 216 EXPERIMENTAL ONE-DIMENSIONAL EXPERIMENTS	
6.2. 6.3. 6.4. PAR CHAF	Selective 6.2.1. 6.2.2. Effects Strong 6.4.1. 6.4.2. 6.4.3. Referent T II. E	Theory / 192 Consequences / 199 of Cross-Correlation / 201 Coupling / 204 A{B} Enhancements / 206 AB{X} Enhancements / 206 Scalar Relaxation / 213 6.4.3.1. Scalar Relaxation of the First Kind / 213 6.4.3.2. Scalar Relaxation of the Second Kind / 216 ces / 217 EXPERIMENTAL ONE-DIMENSIONAL EXPERIMENTS Preparation / 221 Solvent / 221	

7.2.	Setting	Up the Steady-State Difference Experiment / 227				
	7.2.1.	Introduction to the Difference Experiment / 227				
	7.2.2.	Minimizing Subtraction Artifacts / 230				
	7.2.3.	Automatic Multiple Experiments / 234				
	7.2.4.	Irradiation Power and Selectivity / 235				
	7.2.5.	Multiplet Irradiation and SPT Suppression / 240				
	7.2.6.	Timing / 243				
7.3.	Process	sing, Display, and Calculation of Results / 247				
	7.3.1.	General / 247				
	7.3.2.	Reference Deconvolution / 249				
7.4.	Other 1	D Experiments Employing Continuous Saturation / 254				
	7.4.1.	CW Steady-State Integration / 254				
	7.4.2.	The Truncated Driven NOE (TOE) Experiment / 255				
7.5.	Transie	nt Experiments / 258				
	7.5.1.	Selective Pulses / 259				
	7.5.2.	Non-Gradient Transient NOE Experiments / 260				
	7.5.3.	Gradient-Assisted Transient NOE Experiments / 261				
		7.5.3.1. Gradient Selection / 262				
		7.5.3.2. DPFGSE-NOE / 264				
		7.5.3.3. GOESY / 267				
		7.5.3.4. Variations / 272				
		7.5.3.5. Applications and Practicalities / 274				
	7.5.4.	Doubly Selective Experiments / 276				
	Referen	nces / 279				
СНА	PTER 8.	THE TWO-DIMENSIONAL NOESY EXPERIMENT	282			
8.1.	One Di	imension or Two? / 282				
0.11	8.1.1.	The Negative NOE Regime ($\omega \tau_c > 1.12$) / 283				
	8.1.2.	The Positive NOE Regime ($\omega \tau_c < 1.12$) / 283				
8.2.		Principles / 285				
8.3.		ing a NOESY Spectrum / 293				
	8.3.1.	Fixed Delays and Pulse Widths / 294				
	8.3.2.	Acquisition Times t_1 and t_2 and Spectral Widths SW_1 and SW_2 / 296				
	8.3.3.	Quadrature Detection in F_1 and F_2 / 299				
	8.3.4.	Phase-Sensitive NOESY / 304				
8.4.	Phase-C	Cycling, Signal Selection, and Artifact Suppression / 307				
	8.4.1.	Rejection of Nonlongitudinal Contributions During $\tau_{\rm m}$: <i>J</i> -Peak Suppression / 309				

X CONTENTS

	8.4.2.	Other Forms of <i>J</i> -Peaks: Zero Quantum Coherences and Pulse Angle Effects / 310				
	8.4.3.	Axial Peaks / 313				
	8.4.4.	Quadrature Images / 313				
	8.4.5.	t ₁ Noise / 314				
8.5.	Data Pro	ocessing / 316				
	8.5.1.	Zero Filling / 317				
	8.5.2.	Window Functions and Linear Prediction / 318				
	8.5.3.	Symmetrization and t_1 Noise Removal / 320				
	8.5.4.	Integration / 321				
8.6.	Variation	ns / 322				
	8.6.1.	Semi-Selective and Network-Edited Experiments / 322				
	8.6.2.	Other Variants / 327				
	Reference	pes / 328				
CHAR	TED A	OTHER DEVEL ORMENTS	224			
СПАР	TER 9.	OTHER DEVELOPMENTS	331			
9.1.	Heteroni	uclear NOE Enhancements / 331				
	9.1.1.	Non-Specific Heteronuclear NOE Experiments / 331				
	9.1.2.	Specific Heteronuclear NOE Experiments / 335				
9.2.	Editing a	and Spectral Simplification of NOE Experiments / 341				
	9.2.1.	Editing Using the NOE Itself / 341				
	9.2.2.	Editing Using Heteronuclear Scalar Couplings / 346				
		9.2.2.1. X-Filtered NOE Experiments / 347				
		9.2.2.2. X-Separated NOE Experiments / 350				
	9.2.3.	Homonuclear Three-Dimensional NOE Experiments / 359				
	9.2.4.	Editing Using Something Else / 362				
9.3.	Rotating	Frame NOE Experiments / 364				
	9.3.1.	Theory / 365				
		9.3.1.1. Spin-Locking / 365				
		9.3.1.2. Spin-Locked Transverse Dipole-Dipole Relaxation / 367				
		9.3.1.3. Other Effects During Spin-Locking / 371				
	9.3.2.	Practicalities / 379				
9.4.	Manipul	ation of $\omega \tau_{\rm c}$ / 382				
	References / 384					

PAR [®]	T III.	APPLICATIONS	389
		. APPLICATIONS OF THE NOE TO ELUCIDATION	391
10.1.	General	Considerations / 391	
	10.1.1.	Why Structural and Conformational Problems Are the Same / 391	
	10.1.2.	Spectra and Assignments / 393	
	10.1.3.	Reporting Results and Interpretation / 395	
	10.1.4.	Miscellaneous / 397	
10.2.	Aromati	ic Substitution and Ring Fusion Patterns: Simple Cases / 39	98
10.3.	Aromati Cases /	ic Substitution and Ring Fusion Patterns: More Complex 406	
	10.3.1.	Petroporphyrins / 407	
	10.3.2.	Isoquinoline and Related Alkaloids / 409	
10.4.	Double	Bond Isomers / 415	
10.5.	Saturate	ed Ring Systems: Simple Cases / 420	
	10.5.1.	Substituent Stereochemistry / 422	
	10.5.2.	Ring Fusion Stereochemistry / 432	
10.6.	Saturate	ed Ring Systems: More Complex Cases / 442	
	10.6.1.	Pulvomycin / 444	
	10.6.2.	Penitrem A / 447	
	10.6.3.	Other Examples / 450	
	Referen	ces / 453	
		. APPLICATIONS OF THE NOE TO TONAL ANALYSIS	456
11.1.	General	Considerations / 456	
	11.1.1.	Why Structural and Conformational Problems are Different / 456	
	11.1.2.	Multiple Conformations / 458	
11.2.	Local C	Conformational Detail in Small Molecules / 459	
	11.2.1.	Slowly Exchanging Equilibria / 461	
	11.2.2.	Rapidly Exchanging Equilibria: A Hypothetical Example, X -CH ₂ OH $/$ 461	
	11.2.3.	Rapidly Exchanging Equilibria: Real Examples / 463	

11.3. Conformational Analysis of Medium-Sized Molecules / 472

References / 484

OLLA DEED 40	CALCULATING CERLICITIES OF BIODOLVMERC	405
CHAPTER 12.	CALCULATING STRUCTURES OF BIOPOLYMERS	483

12.1.	Introduction	/ 485
14.1.	muoduction	/ +02

- 12.2. Restraints / 486
 - 12.2.1. Assigning NOE Restraints / 487
 - 12.2.1.1. Using Only (¹H, ¹H) NOESY Data / 488
 - 12.2.1.2. Using Preliminary Structural Data / 489
 - 12.2.1.3. Using Heteronuclear Labeling / 490
 - 12.2.2. Measuring NOE Restraints / 491
 - 12.2.3. Calibrating NOE Restraints / 493
 - 12.2.4. Averaging in Equivalent Groups / 499
 - 12.2.4.1. Pseudoatom Corrections / 500
 - 12.2.4.2. Multiplicity Corrections / 503
 - 12.2.4.3. r^{-6} Summation / 506
 - 12.2.5. Stereoassignments and Torsion Angle Restraints / 506
 - 12.2.6. Other Types of Restraints / 511
- 12.3. Calculating Structures / 515
 - 12.3.1. Distance Geometry Calculations / 517
 - 12.3.1.1. The Exact Case / 518
 - 12.3.1.2. Distance Geometry Applied to NMR Structure Determination / 519
 - 12.3.1.3. Distance Geometry in Torsion Angle Space / 522
 - 12.3.2. Restrained Molecular Dynamics Calculations / 523
 - 12.3.3. Simulated Annealing Calculations / 527
 - 12.3.4. Other Methods / 529
- 12.4. Assessing and Describing NMR Structures / 530
 - 12.4.1. Global Precision: Overall Root Mean Squared Deviations / 531
 - 12.4.2. Local Precision: Local RMSD and Angular Order Parameters / 533
 - 12.4.3. Assessing the Quality of Structures / 536
- 12.5. Refinement / 540
 - 12.5.1. General / 541
 - 12.5.2. Specific Protocols for Refinement / 543

References / 546

CHAPTER 13. BIOPOLYMERS

550

- 13.1. Peptides and Proteins / 550
 - 13.1.1. Assignment: Heteronuclear Methods / 554

	13.1.2.	Assignment: (¹ H, ¹ H) NOE-Based Methods / 555		
	13.1.3.	Structure Determination of Protein Monomers / 556		
		13.1.3.1. Small Rigid Unlabeled Proteins / 556		
		13.1.3.2. Larger Rigid Labeled Proteins / 558		
		13.1.3.3. Conformationally Mobile Proteins / 560		
	13.1.4.	Structure Determination of Symmetric Protein		
		Oligomers / 562		
	13.1.5.			
		Experiments with Proteins / 563		
13.2.	•	leotides / 566		
		Structures and Conformations / 567		
	13.2.2.	Assignment / 574		
		13.2.2.1. Duplex DNA / 574		
		13.2.2.2. RNA / 576		
		13.2.2.3. Other Nucleotides / 577		
	13.2.3.	Structure Calculation / 577		
		13.2.3.1. Sequence Dependent Conformation in		
		Duplexes / 577		
		13.2.3.2. Non-Helical Conformations / 581		
13.3.	•	ccharides / 583		
	13.3.1.			
		Conformation / 585		
13.4.	-	xes / 587		
		Drug-Protein Complexes / 589		
		Drug-DNA Complexes / 589		
	13.4.3.	Protein-Nucleic Acid Complexes / 590		
	Referen	ces / 592		
		EQUATIONS FOR ENHANCEMENTS INVOLVING	FOC	
GRU	J29 0F	EQUIVALENT SPINS	596	
APPENDIX II. QUANTUM MECHANICS AND TRANSITION				
PROE	BABILITI	E5	599	
INDEX			611	

PREFACE

When we completed the first edition of this book, just over 10 years ago, we did not expect to undertake rewriting it for a long time to come. It is certainly true that the theory behind the NOE has altered little in the last 10 years (as reflected by the fact that Chapters 1, 2, 3, and 6 are largely unchanged). However, there has continued to be a very rapid growth in the range of applications, particularly in the macromolecular area.

To a large extent, this has dictated the changes we have made to the book. The biggest change from the first edition is the inclusion of a completely new chapter (Chapter 12), which deals with the way NOE enhancements are used to calculate structures of biomolecules. Chapter 13, which discusses applications of the NOE to calculations of macromolecular structures, has been largely rewritten. Much the same is true for Chapter 4, which deals with the kinetics of the NOE; although the theory is largely the same, we have now placed much more emphasis on transient NOE experiments (such as 2D and gradient-selected experiments), which have seen a major increase in use over the last 10 years.

Interest in internal motions in macromolecules was the main driving force behind the changes we have made to Chapter 5 (which also includes a new section on NOE enhancements involving water, and a major update on the transferred NOE), and a large number of new approaches have been incorporated into Chapters 7 (particularly gradient enhanced experiments) and 8 (network editing). Chapter 9 ("Other Developments") remains the home for "miscellany" and contains a lot of new material, much of it concerned with the new experiments that became available in the wake of isotope labeling.

We have not significantly changed the content of the chapters dealing with applications to small molecules (Chapters 10 and 11), partly because neither

of us is still active in that field, but more importantly because we feel the examples we gave in the first edition still represent a balanced selection that demonstrates coherently what is usefully possible. It is true that some types of structural problem are now more frequently solved by other means (e.g., positional isomers on aromatic rings might now be more frequently distinguished using long-range ¹³C-¹H correlation than by using the NOE), but we felt illustration of the utility of such NOE experiments was still important.

Overall, we have thus broadened the scope of the first edition somewhat, feeling that to exclude (for example) applications of the NOE to relaxation measurements, because it does not relate directly to structure or conformation, was unnecessarily restrictive.

One of the aims we set ourselves in writing this second edition was to make the book shorter. The most casual glance will reveal that we have achieved the opposite of success in this regard. The reason, unfortunately, is simple: as George Bernard Shaw (or was it Blaise Pascal?) is supposed to have said, "I am sorry to have written you a long letter, but I did not have time to write you a short one." For this and other reasons, we feel it is safe to predict that there will be no third edition.

David Neuhaus Michael P. Williamson

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PREFACE TO THE FIRST EDITION

In the past ten years a quiet revolution in the applications of NMR to organic chemistry and biochemistry has occurred. One of the main changes has been a great increase in the use of the nuclear Overhauser effect (NOE) to solve structural and conformational problems, which in turn was largely due to experimental developments such as the advent of NOE difference spectroscopy and the two-dimensional NOE experiment (NOESY). In this book we have tried to bring chemists and biochemists who are not NMR specialists up to date with these trends. The book aims to provide readers with sufficient background information to enable them to apply the NOE successfully within their own work, and also assess critically other papers in which the NOE is used to solve structural problems.

Because the NOE is transmitted directly through space, it is uniquely well suited to revealing the spatial arrangements of nuclei within molecules, but its interpretation is particularly vulnerable to the overapplication of inappropriate equations. For this reason, we have tried above all to emphasize a clear understanding of the underlying concepts, since only through an awareness of these can the most useful NOE experiments be devised and interpreted. Unavoidably, we have had to use a certain amount of mathematics to do this, but it is kept to a minimum, consistent with a clear explanation. Perhaps contrary to first appearances, most of the expressions in this book require no more than elementary mathematical skills.

As the title implies, our treatment is deliberately restricted to those aspects of the NOE that are relevant to the study of stereochemistry and conformation; topics such as dipolar interactions in solids, and the use of the NOE to study relaxation phenomena, kinetics, or exchange, have therefore been omitted.

XX PREFACE

The book is organized into three parts, dealing with theory, experimental practice, and an illustrative overview of applications of the NOE to problems in chemistry and biochemistry.

SI units have been used throughout, as far as possible. This results in occasional numerical factors in some expressions (e.g., the factor $\mu_0/4\pi$, which appears in expressions relating relaxation rates to internuclear distances).

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SYMBOLS, ABBREVIATIONS, AND UNITS

Throughout, subscript letters are used to denote spin states and transitions; thus, N_{α} is the number of molecules in spin state α , and R_I is the spin-lattice relaxation rate of spin I. Superscript letters are used to denote different molecules or conformations; thus, S^a is the signal from spin S in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b is the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b in molecules with conformation A^b in the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b in the spin-lattice relaxation rate of spin A^b in molecules with conformation A^b in the spin-lattice relaxation rate of spin A^b in the spin A^b is the spin-lattice relaxation rate of spin A^b in the spin A^b is the spin A^b in the spin A^b in the spin A^b is the spin A^b in the spin A^b in the spin A^b is the spin A^b in the spin A^b in the spin A^b in the spin A^b is the spin A^b in the spin A^b

analog-to-digital converter
acquisition time
field strength in Hz
field strength in rad s ⁻¹
applied magnetic field (tesla)
irradiating magnetic field (tesla)
continuous wave
distance between proton A in residue i and proton B in residue
j in a peptide or protein
dwell time in t_2 (s)
one-dimensional
two-dimensional
three-dimensional
decibels
distance geometry
double pulsed field gradient spin echo
double-quantum filtered 2D correlation spectroscopy
fractional NOE enhancement of I on saturating S

free induction decay FID FT Fourier transform

 $g(\tau)$ correlation function, usually $\exp(-\tau/\tau_c)$ **GOESY** gradient enhanced NOE spectroscopy Planck's constant = 6.626×10^{-34} Js h

Planck's constant divided by $2\pi = 1.055 \times 10^{-34} \text{ Js}$ ħ

HMOC heteronuclear multiple quantum correlation (also coherence) **HSQC** heteronuclear single quantum correlation (also coherence) Ι nuclear spin quantum number, or general symbol for a nucleus; in expressions concerning NOE experiments, I generally repre-

sents the nucleus at which the enhancement is measured.

longitudinal component of I magnetization

 $I_z \\ I_z^0$ equilibrium value of I_z

IN t_1 increment in a 2D experiment (s) **ISPA** isolated spin pair approximation spin-spin coupling constant (Hz) J

spectral density function, usually $2\tau_c/(1 + \omega^2 \tau_c^2)$ $J(\omega)$ Boltzmann's constant = $1.38 \times 10^{-23} \text{ J K}^{-1}$ k macroscopic magnetization vector (J T⁻¹) \mathbf{M}

 M_{τ} z component of M

 M_{xy} transverse component of M MAS magic angle spinning

number of molecules (population) in state i N_i

population deviation from equilibrium of state i, namely N_i – n_i

 N_i^0

NOE nuclear Overhauser effect

NOESY nuclear Overhauser effect correlation spectroscopy

 $R_I^{\rm DD}$ dipolar contribution to the spin-lattice relaxation rate of I, defined as $W_{0IS} + 2W_{1I} + W_{2IS} + \sum_{x} (W_{0IX} + W_{2IX})$ (cf. Eqs. 2.31)

and 3.8) (s^{-1})

longitudinal relaxation rate of $I = R_I^{DD} + \rho_I^*$, roughly equal to $R_I \equiv R_{1I}$ the inverse of the selective T_1 value (cf. Section 2.3.1) (s⁻¹)

 R_2 transverse relaxation rate (s⁻¹)

internuclear distance

effective internuclear distance sensed by the NOE in the pres $r_{\rm Tropp}$ ence of rapid internal motion; most commonly used to denote the effective distance to a methyl group (cf. Section 5.5.2)

RF radiofrequency

root mean squared deviation (or difference); usually implies root rmsd

mean squared atomic deviation when comparing structures

ROE rotating frame NOE; also known as CAMELSPIN

ROESY rotating frame 2D NOE spectroscopy

S general symbol for a nucleus; in expressions concerning NOE experiments, S generally represents the nucleus which is satu-

rated (S can also represent the electron spin quantum number).

Also, generalized order parameter for describing the geometrical extent of internal motions relative to the molecular frame (cf.

Section 5.4)

 S^2 square of the generalised order parameter (cf. Section 5.4)

 S^{ang} angular order parameter (cf. Section 12.4.2) S_{z} longitudinal component of S magnetization

 S_z^0 equilibrium value of S_z SPT selective population transfer

spectral width in F_1 of a multidimensional experiment (Hz) SW_1 SW_2 spectral width in F_2 of a multidimensional experiment (Hz) SW_3 spectral width in F_3 of a multidimensional experiment (Hz) incremented time in 2D experiments; first incremented time in t_1

multidimensional experiments

maximum time reached by t_1 in the last increment of a 2D or $t_{1 \, \text{max}}$

multidimensional experiment (s)

detection period in 2D experiments; second incremented time in t_2

multidimensional experiments

maximum time reached by t_2 (s) $t_{2\,\mathrm{max}}$ t_3 detection period in 3D experiments maximum time reached by t_3 (s) $t_{3\,\mathrm{max}}$

Ttemperature (K)

normally used to mean $T_1^{\text{nonselective}}$ (see below) T_1

 $T_1^{
m nonselective}$ nonselective longitudinal relaxation time, obtained by following the course of longitudinal relaxation after a nonselective pulse; approximately given by $(R_I + \sigma_{IS} + \Sigma_X \sigma_{IX})^{-1}$ (s) (cf. Section

 $T_1^{
m selective}$ selective longitudinal relaxation time, obtained by following the course of longitudinal relaxation after a selective pulse; approx-

imately equal to R_I^{-1} (s)

 $T_{1\rho}$ spin-lattice relaxation time in the rotating frame (s)

 T_2 spin-spin, or transverse, relaxation time (s) $T_{2\rho}$ spin-spin relaxation time in the rotating frame (s)

 T_2^* decay constant for free precession, which is shorter than T_2 be-

cause of inhomogeneous broadening (s)

 T_{2o}^* decay constant for free precession in the rotating frame (s)

 $t_{
m D}$ relaxation delay (s) **TMS** tetramethylsilane

TOCSY total correlation spectroscopy

truncated driven NOE TOE

TPPI time-proportional phase incrementation

TRNOE transferred NOE **VDW** van der Waals VT variable temperature transition probability (s⁻¹) W

X	general symbol for a nucleus; in expressions concerning NOE experiments in this book, X represents all nuclei other than I and S (when not italicized, X can also refer to a non-hydrogen
	nucleus)
X_z	longitudinal component of X magnetization
X_z^0	equilibrium value of X_z
α, β	spin states for a spin 1/2 nucleus
γ	gyromagnetic ratio (rad T ⁻¹ s ⁻¹)
η	viscosity (cP)
$oldsymbol{\eta}_{ ext{max}}$	maximum enhancement attainable in a two-spin system at a
	given value of $ au_{\rm c}$ and ω
μ	magnetic dipole moment (J T ⁻¹)
$oldsymbol{\mu}_0$	permeability constant in a vacuum = $4\pi \times 10^{-7}$ kg m s ⁻² A ⁻²
ν	precession rate (Hz)
$ u_0$	reference frequency of rotating frame (Hz)
$ ho_{\scriptscriptstyle IS}$	direct dipole-dipole relaxation rate between I and S , defined by
	$\rho_{IS} = W_{0IS} + 2W_{1I} + W_{2IS}$ [Eq. 2.30] (s ⁻¹)
$ ho^*$	external spin-lattice relaxation rate (s ⁻¹)
$oldsymbol{ ho}_1$	alternative notation for ρ_{IS} , useful when comparing the NOE with
	the ROE
$ ho_2$	direct relaxation rate in the rotating frame (s ⁻¹)
$\sigma_{\scriptscriptstyle I}$	chemical shift tensor for nucleus I (cf. Sections 5.4 and 6.3)
$\sigma_{\scriptscriptstyle IS}$	cross-relaxation rate between I and S , defined by $\sigma_{IS} = W_{2IS} - W_{0IS}$ (Eq. 2.29) (s ⁻¹)
$\sigma_{\scriptscriptstyle 1}$	alternative notation for σ_{IS} , useful when comparing the NOE
	with the ROE
σ_{2}	cross-relaxation rate in the rotating frame (s ⁻¹)
$ au_{ m c}$	rotational correlation time (for overall molecular tumbling) (s)
$ au_{ m e}$	correlation time for an internal motion (cf. Section 5.4) (s)
$ au_{ m m}$	mixing time in a 2D or multidimensional NOE experiment (s)
au	buildup time in kinetic NOE experiment; also general symbol
	for a fixed delay other than $\tau_{\rm m}$ (s)
ω_0	Larmor precession frequency due to \mathbf{B}_0 (rad s ⁻¹)
ω_1	Larmor precession frequency due to \mathbf{B}_1 (rad s ⁻¹)
{ }	denotes irradiated spin
< >	denotes average or expectation value or conformational average
\langle , \rangle	Dirac notation for a wave function (ket) and its complex con-
	jugate (bra), respectively

INTRODUCTION

The fundamentals of the NOE were described very early in the history of NMR, in a classic paper by Solomon published in 1955. This paper included the first experimental demonstration of the NOE. Solomon's work followed studies on nuclear spin relaxation by Bloembergen, Purcell, and Pound, and Overhauser's original prediction that saturation of electrons in a metal would produce a large polarization of the metal nuclear spins. 4

Solomon's work then lay almost dormant for some years until the advent of double resonance techniques^{5,6} led to the more widespread availability of spectrometers having a decoupler. Papers appeared by Kaiser on distance estimation using the intermolecular ¹H{¹H} NOE,⁷ and by Lauterbur on signal enhancement in ¹³C spectra.⁸

The first paper to demonstrate the power of the NOE in structural problems was by Anet and Bourn in 1965. The next few years saw a surge in the number of papers reporting applications of the NOE, especially after publication of the book by Noggle and Schirmer. On the CW spectrometers of the day, these experiments involved comparison of integrals to measure steady-state NOE enhancements.

Noggle and Schirmer's book surveyed the field as it stood in 1971, and for some years there was little that needed to be added. There were two major advances in the decade or so following 1971. The first was the application of the NOE to large molecules. In 1972, Balaram et al. 11 observed a negative NOE enhancement between a protein and a peptide binding to it, and they showed that negative enhancements are a general property of large molecules. Further applications to large molecules soon revealed the problem of spin diffusion, and the truncated driven NOE (TOE) experiment was suggested by Wagner and Wüthrich 12 as a means of getting around this difficulty.

The second major advance was the introduction of the two-dimensional NOE experiment, largely by Ernst's group. The first use of the experiment to measure NOE enhancements (as opposed to chemical exchange) was published in 1980, ^{13,14} and the experiment has come to play a central role in work with biological macromolecules. This area has burgeoned still further with the arrival in the late 1980s of isotopically labeled proteins and the many multidimensional NOE experiments that can be applied to them.

Both these advances could not have had the impact they did without developments in instrumental techniques. The first FT NMR spectrometers became available around 1970, and brought much improved sensitivity. To observe small NOE enhancements, one needs not only sensitivity but also instrumental stability, and this was provided by superconducting magnets, which arrived at the end of the 1970s. The third in this trio of instrumental advances was the advent of dedicated computers, which could control the cycling of decoupler frequencies and gating. A direct result of these changes was the NOE difference experiment, which emerged at about this time. A much more recent development is the use of pulsed field gradients, which have found applications in almost all parts of NMR. Gradients provide little benefit for 2D NOESY experiments, but promise to revolutionize 1D measurements.¹⁵

Although the theory of the NOE as applied to organic molecules has scarcely changed since 1971, its application has increased dramatically. Using the techniques available in 1971, enhancements smaller than 5% could be seen only with difficulty. Now it is not particularly difficult to detect enhancements of less than 1%, or even 0.1% using newer gradient assisted experiments. Hall and Sanders¹⁶ point out that the maximum possible size of NOE enhancement seen between 1,3 diaxial protons in a cyclohexane chair is about 3%. This size of enhancement was just below the threshold of detection in the CW era, but it is well within modern instrumental capabilities.

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